

Time's arrow at the nanoscale

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Macroscopic and microscopic machines

steam engine





What do the laws of thermodynamics "look like" at the nanoscale?

Work and free energy: a macroscopic example ...



Irreversible process:

- 1. Begin in equilibrium
- Stretch the rubber band
 W = work performed
- 3. End in equilibrium

 $\lambda = A$

 $\lambda : A \rightarrow B$

 $\lambda = B$

Work and free energy: a macroscopic example ...



Clausius inequality :

$$W_{A \to B} \ge \Delta F \equiv F_B - F_A$$

$$\left(\int_{A}^{B} \frac{dQ}{T} \leq \Delta S\right)$$

... and a microscopic analogue





- 1. Begin in equilibrium
- Stretch the molecule
 W = work performed
 - $\lambda : A \rightarrow B$

 $\lambda = A$

 $\lambda = B$

- 3. End in equilibrium
- 4. Repeat

... fluctuations are now important !



So what's new?

Fluctuations in W satisfy strong and unexpected laws.

e.g.
$$\left\langle e^{-\beta W} \right\rangle = e^{-\beta \Delta F}$$
 C.J., PRL 1997 $\left(\beta = \frac{1}{k_B T}\right)$

... places a strong constraint on the work distribution





 $-\mathsf{W}_\mathsf{R} \leq \Delta\mathsf{F} \leq \mathsf{W}_\mathsf{F}$

Kelvin-Planck statement of 2nd Law: $W_F + W_R \ge 0$

(no free lunch)



Kelvin-Planck statement of 2nd Law: $\langle W \rangle_F + \langle W \rangle_R \ge 0$ (no free lunch... on average)



$$\frac{\rho_F(+W)}{\rho_R(-W)} = \exp[\beta(W - \Delta F)]$$

G.E. Crooks, PRE 1999

Unfolding & refolding of ribosomal RNA

$$\frac{\rho_{unfold}(+W)}{\rho_{refold}(-W)} = \exp[\beta(W - \Delta F)]$$



Nonequilibrium work relations

macro	micro
$W \ge \Delta F$	$\langle W \rangle \ge \Delta F$
$-W_R \leq \Delta F \leq W_F$	$-\left\langle W\right\rangle _{R}\leq\Delta F\leq\left\langle W\right\rangle _{F}$

Nonequilibrium work relations



closely related to *Fluctuation Theorems* for entropy production* and to early work by Bochkov & Kuzovlev[†]

> * Evans, Cohen, Morriss, Gallavotti, Searles, Kurchan, Lebowitz, Spohn, Maes + many others
> † JETP 1977, 1979

Nonequilibrium work relations

macro	micro
$W \ge \Delta F$	$\left\langle e^{-\beta W} \right\rangle = e^{-\beta \Delta F}$
$-W_R \leq \Delta F \leq W_F$	$\frac{\rho_F(+W)}{\rho_R(-W)} = \exp[\beta(W - \Delta F)]$
	•••

- apply far from equilibrium
- validated in experiments
- applications: analysis of single-molecule experiments
 & computational thermodynamics
- irreversibility and the arrow of time at the nanoscale

The thermodynamic arrow of time

Make a movie of a thermodynamic process $(A \rightarrow B)$



Macroscopic system: $W > \Delta F$ if the movie is run forward $W < \Delta F$ if the movie is run backward

... perfect correlation between sign(W- Δ F) and the chronological ordering of events.

The arrow of time is *sharp*.

The thermodynamic arrow of time

Make a movie of a thermodynamic process $(A \rightarrow B)$



Microscopic system:

typically $W > \Delta F$ if the movie is run forwardtypically $W < \Delta F$ if the movie is run backward

... but there are exceptions (fluctuations).

The arrow of time is *blurred*.

How frequently do these exceptions occur?

Irreversibility in microscopic systems



 $\rightarrow \rho(W)$ is *exponentially suppressed* in the thermodynamically forbidden region



<u>Hysteresis</u>

The system evolves via one sequence of states during the forward process (stretching), but follows a different path during the reverse process (contraction).



Hysteresis = ?



The probability to observe one sequence of events (γ_F) during the forward process is different from that of observing the conjugate sequence (γ_R) during the reverse process.

$$P_F[\gamma_F] \neq P_R[\gamma_R]$$



Guessing the direction of time's arrow

You are shown a movie depicting a thermodynamic process, A→B. Task: determine whether you are viewing the events in the order in which they actually occurred, or a movie run backward of the reverse process.

e.g.







Relative entropy and time's arrow

Relative entropy D[f | g] provides a measure of the difference between two probability distributions f and g.

$$D[f \mid g] = \int f \ln \frac{f}{g} \ge 0$$

 $\begin{array}{lll} \mbox{forward process } \lambda : A \rightarrow B & p_F(x,t) \\ \mbox{reverse process } \lambda : A \leftarrow B & p_R(x,t) \end{array}$

use relative entropy to quantify irreversibility, by comparing p_F with p_R

notation: $W_{diss} = W - \Delta F$, dissipated work

Reversible processes λ=A $\lambda = \lambda_{int}$ **λ=B** p_F(x,t) p_F(x,t) p_F(x,t) t=0 $t=\tau/2$ **t**= τ t= τ/2 t= 0 t=τ p_R(x,t) p_R(x,t) $p_R(x,t)$

 $p_F(x,t) = p^{eq}(x; \lambda^F(t)) = p_R(x,\tau-t)$

Irreversible processes



 $p_F(x,t) \neq p^{eq}(x; \lambda^F(t)) \neq p_R(x,\tau-t)$



use relative entropy to quantify both hysteresis & lag:

 $D\left[p_{F} \mid p_{R}\right]_{\lambda} \leq \beta \langle W_{diss} \rangle_{F} \quad \text{Kawai, Parrondo, Van den Broeck, PRL 2007}$ $D\left[p_{F} \mid p^{eq}\right]_{t} \leq \beta \langle W_{diss}(t) \rangle_{F} \quad \text{Vaikuntanathan \& CJ, EPL 2009}$



$$D[p \mid p^{eq}]_{t} \leq \beta \langle W_{diss}(t) \rangle$$

$$\rightarrow D[p_{\tau} \mid p_{B}^{eq}] \leq \beta \langle W_{diss} \rangle$$

$$\rightarrow \langle W \rangle \geq \Delta F + k_{B} T D[p_{\tau} \mid p_{B}^{eq}]$$

2nd Law



At the nanoscale ...

• the 2nd law can be represented by equalities

$$\left\langle e^{-\beta W} \right\rangle = e^{-\beta \Delta F}$$
, $\frac{\rho_F(+W)}{\rho_R(-W)} = \exp[\beta(W - \Delta F)]$, & others

• the arrow of time is blurred, but can be quantified